

# DISCOVERY OF AN $\sim 23 M_{\text{Jup}}$ BROWN DWARF ORBITING $\sim 700$ AU FROM THE MASSIVE STAR HIP 78530 IN UPPER SCORPIUS

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## ABSTRACT

We present the discovery of a substellar companion on a wide orbit around the  $\sim 2.5 M_{\odot}$  star HIP 78530, which is a member of the 5 Myr old Upper Scorpius association. We have obtained follow-up imaging over two years and show that the companion and primary share common proper motion. We have also obtained *JHK* spectroscopy of the companion and confirm its low surface gravity, in accordance with the young age of the system. A comparison with DRIFT-PHOENIX synthetic spectra indicates an effective temperature of  $2800 \pm 200$  K and a comparison with template spectra of young and old dwarfs indicates a spectral type of  $M8 \pm 1$ . The mass of the companion is estimated to be  $19\text{--}26 M_{\text{Jup}}$  based on its bolometric luminosity and the predictions of evolutionary models. The angular separation of the companion is  $4''.5$ , which at the distance of the primary star, 156.7 pc, corresponds to a projected separation of  $\sim 710$  AU. This companion features one of the lowest mass ratios ( $\sim 0.009$ ) of any known companion at separations greater than 100 AU.

*Key words:* brown dwarfs – planetary systems – stars: formation – stars: low-mass – stars: pre-main sequence

*Online-only material:* color figures

## 1. INTRODUCTION

The recent direct imaging discoveries of very low mass substellar companions to stars—especially those orbiting the young stars IRXS J160929.1–210524 (Lafrenière et al. 2008, 2010), HR 8799 (Marois et al. 2008), Fomalhaut (Kalas et al. 2008), AB Pic (Chauvin et al. 2005), HN Peg (Luhman et al. 2007), and CT Cha (Schmidt et al. 2008)—have raised new questions about the formation mechanisms of planets and brown dwarfs. These companions have separations of tens to several hundred AU and mass ratios  $\lesssim 0.02$  relative to their primaries. Several theoretical ideas are under discussion to explain such properties, but they all run into some difficulties. Core accretion models favor formation of giant planets close to the “snow line,” i.e., at  $< 10$  AU (Pollack et al. 1996). Gravitational instability could produce massive planets at large radii, but would require unusually large and massive circumstellar disks (e.g., Vorobyov & Basu 2010). Fragmentation of pre-stellar cores during collapse does not lead to such extreme binary systems easily (e.g., Bate et al. 2003). Thus, some theorists have recently proposed that dynamical instabilities within planetary systems that originally formed multiple giant planets could scatter some of them to large separations (Scharf & Menou 2009; Veras et al. 2009). However, the more massive planets are less likely to reach the widest orbits.

Here, we report on the direct imaging discovery, common proper motion confirmation, and multi-band spectroscopy of an  $\sim 23 M_{\text{Jup}}$  companion seen  $\sim 700$  AU from HIP 78530, which is an  $\sim 2.5 M_{\odot}$  star (spectral-type B9V) in the Upper Scorpius young association. The separation of this new companion places it among the widest known substellar companions to stars and its extreme mass ratio—among the lowest currently known—is comparable to those of directly imaged planets, even though its mass is well above the deuterium-burning threshold. Together

with the companions mentioned above, this new substellar companion presents a good challenge to all formation scenarios and contributes to blurring the distinction between giant planets and brown dwarfs even further.

## 2. OBSERVATIONS AND DATA REDUCTION

### 2.1. Imaging

The discovery presented in this paper was made as part of a direct imaging search for new stellar and substellar companions around about 90 stars in the Upper Scorpius region. The overall target sample was built by randomly selecting, from the list of Upper Scorpius stars in Carpenter et al. (2006), an equal number of stars in each of five equal logarithmic mass bins over the range  $\sim 0.15\text{--}5 M_{\odot}$ ; the spectral types range from B0 to M5. The observations were made using the NIRI camera (Hodapp et al. 2003) and the ALTAIR adaptive optics system (Herriot et al. 2000) at the Gemini North Telescope. Except for a few faint targets, the target stars themselves were used for wave-front sensing. The field lens of ALTAIR was used to reduce the effect of anisoplanatism. The first epoch imaging of HIP 78530 was done on 2008 May 24 in the narrowband filter  $K_{\text{continuum}}$  centered at  $2.0975 \mu\text{m}$ . For sky subtraction, we used five dither positions corresponding to the corner and center of a square of side  $10''$ . At each position, we obtained one co-addition of 12 0.5 s integrations in fast, high read-noise mode, followed by one single 10 s integration in slow, low read-noise mode. At each position, this provides an unsaturated image of the target star and a much deeper image of the field that can be readily spatially registered and scaled in flux. The full width at half-maximum for this data set is  $0''.075$  and the Strehl ratio is 0.34.

The initial image of HIP 78530 revealed an interesting faint nearby source which, based on a comparison with observations made by Kouwenhoven et al. (2005, 2007), was very likely

to be a true co-moving companion. To confirm the common proper motion of this candidate and get additional photometry measurements, follow-up imaging observations in  $J$ ,  $H$ , and  $K'$  were obtained on 2009 July 2 using the same instrument as before. To avoid saturation in these broad filters, we used a  $512 \times 512$  subarray and kept the strategy of a short, unsaturated image immediately followed by a longer, saturated image at each of five dither positions. For each wavelength, three sky frames were obtained at offsets of  $>15''$ . The short exposure times were 0.055 s with 40 co-additions, while the long exposures were 6 s, 5 s, and 4 s with one co-addition in  $J$ ,  $H$ , and  $K'$ , respectively. To increase the significance of our common proper motion confirmation further, another follow-up observation was made on 2010 August 30 using the  $K_{\text{cont}}^{2.09}$  filter, full frame, and one co-addition of 12 0.6 s integrations followed by a single 10 s integration at each of five dither positions. Other follow-up data were acquired in Spring and Summer 2010 but owing to an overseen difference in the instrumental setup, these data suffer from large systematic astrometric errors and are not used here.

The imaging data were reduced using custom IDL routines. For the 2008 and 2010 imaging data, a sky frame was constructed by taking the median of the images at all dither positions after masking out the regions dominated by the target's signal, while for the 2009 imaging data, the median of the three sky frames was obtained. After subtraction of this sky frame, the images were divided by a normalized flat field. Then, isolated bad pixels were replaced by the interpolated value of a third-order polynomial surface fit to the good pixels in a  $7 \times 7$  pixels box, while clustered bad pixels were simply masked out. Next, the images were distortion corrected using the distortion solution provided by the Gemini staff.<sup>5</sup> Finally, the long-exposure images were properly scaled in intensity and merged with their corresponding short-exposure images. The color composite  $JHK$  image from the 2009 follow-up observations is shown in Figure 1.

## 2.2. Spectroscopy

The likely companion detected in the imaging data was very faint and, assuming it was a true companion, its magnitude suggested a substellar mass. Thus, to verify its late-type nature, we obtained spectroscopic observations of it in  $H$ ,  $K$ , and  $J$  on 2009 July 2, 2009 July 3, and 2009 August 8, respectively, using the integral field spectrograph NIFS (McGregor et al. 2003) with the ALTAIR adaptive optics system at the Gemini North Telescope. The  $J$  grating was used with the  $ZJ$  blocking filter, the  $H$  grating with the  $JH$  blocking filter, and the  $K$  grating with the  $HK$  blocking filter. The spectral resolving power is  $\sim 6000$  in  $J$  and  $\sim 5300$  in both  $H$  and  $K$ . Given the large angular separation of the companion ( $\sim 4''.5$ ), the primary star is not visible in the  $3'' \times 3''$  NIFS field of view. In each band, the companion was positioned near the center of the field of view and was dithered by  $0''.7$  along a line between each of five exposures to enable good sky subtraction. The individual exposure times were 480 s, 300 s, and 240 s in  $J$ ,  $H$ , and  $K$ , respectively, in low read-noise mode. After the  $J$  and  $H$  sequences, the A0V star HIP 79229 was observed at a similar air mass for telluric and instrumental transmission correction, while the A0V star HIP 73820 was observed just before the  $K$ -band sequence for the same purpose.

<sup>5</sup> The distortion correction used is centrosymmetric and is given by  $r = r' + 1.32 \times 10^{-5} r'^2$ , where  $r$  and  $r'$  are respectively the distortion-corrected and distorted radial pixel distances from the array center.

**Table 1**  
Astrometric Measurements

Epoch	Band	$\rho$ (")	P.A. (deg)
2008.3940	$K_{\text{cont}}^{2.09}$	$4.529 \pm 0.006$	$140.32 \pm 0.10$
2009.4998	$JHK'$	$4.533 \pm 0.006$	$140.35 \pm 0.10$
2010.6605	$K_{\text{cont}}^{2.09}$	$4.536 \pm 0.006$	$140.27 \pm 0.10$

The reduction of the NIFS data, up to the reconstruction of the data cube, was made using the Gemini IRAF pipeline. The steps covered in this pipeline are sky background subtraction, flat-field and bad pixel correction, spatial and spectral calibration, and data cube reconstruction. The spatial sampling of the reconstructed data cube is  $0''.043 \text{ pixel}^{-1}$ . The instrumental/telluric transmission correction and spectrum extraction, which could be done using the Gemini pipeline, were instead done using custom IDL.

The center of each point-spread function (PSF) was first registered to a common position in all spectral slices and all cubes of the sequence; the center positions were calculated by fitting a two-dimensional Gaussian function. The spectrum of the source was extracted by summing the flux in a circular aperture of diameter 5 pixels. The spectrum of the telluric standard was corrected for its spectral slope using a 9520 K blackbody curve, and any hydrogen line absorption was removed by dividing out a Voigt profile fit. The companion spectrum was divided by the standard spectrum to correct for telluric and instrumental transmission. Finally, the median of the five spectra was obtained.

## 3. ANALYSIS AND RESULTS

A summary of the main properties of HIP 78530 and its companion, taken from the literature or derived in this section, is presented in Table 2.

### 3.1. Common Proper Motion

The relative position of the companion and primary was determined by fitting a two-dimensional Gaussian function to each component. The pixel positions were converted to arcseconds using the 21.4 mas pixel scale of NIRS with ALTAIR and the field lens as indicated on the instrument Web site;<sup>6</sup> the position angle (P.A.) of the image was taken from the FITS header. The astrometric errors, 6 mas in separation and  $0''.1$  in P.A., were estimated in Lafrenière et al. (2010) using the position of background stars for a similar sequence of observations; they are mostly dominated by residual distortion errors. We also note that the pixel scale of ALTAIR/NIRS with the field lens has not been extensively calibrated and that our values may be systematically different from measurements made using other instruments or telescopes, but they should be internally consistent. The astrometry measurements are given in Table 1 and shown in Figure 2 in comparison with the changes expected over time for a distant background star, based on the proper motion and distance of the primary star (taken from van Leeuwen 2007). Our measurements are consistent with common proper motion but inconsistent, at a level of  $\sim 6\sigma$ , with the changes expected for a distant, motionless background star. This clearly indicates that the companion is co-moving with the primary.

The accuracy of our astrometry measurements can be verified using a faint background star detected both in 2008 and

<sup>6</sup> <http://www.gemini.edu/sciops/instruments/altair/field-lens-option>



**Figure 1.** Color composite *JHK* image of HIP 78530 and its substellar companion (bottom left) obtained with NIRI/ALTAIR at the Gemini North Telescope. The intensity scaling of each color is proportional to the incident photon flux in equal fractional bandwidths.

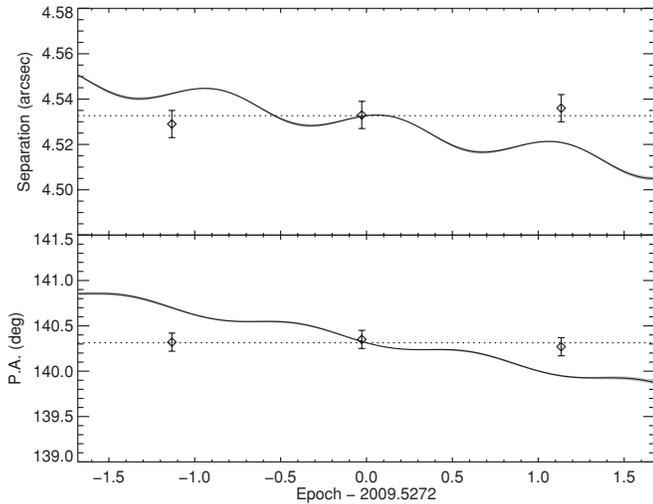
2010. This source,  $\sim 10.1$  mag fainter than the primary, had a separation of  $9''.451$  and P.A. of  $43^\circ 49'$  at epoch 2008.3940, and  $9''.493$  and  $43^\circ 48'$ , respectively, at epoch 2010.6605. These values, showing a  $\sim 7\sigma$  change in separation and  $\sim 1\sigma$  in P.A., are fully consistent with expectations for a distant, motionless background star.

The new companion we report in this paper was detected previously by Kouwenhoven et al. (2005, 2007) in data obtained in 2000 and 2001 with ADONIS at the ESO 3.6 m telescope. They report a separation of  $4''.54 \pm 0''.01$  and P.A. of  $139^\circ 7' \pm 0'.3$ . They note that they were not able to conclude whether this object was a bound companion or a background object based on *HK* photometry, and they have not followed up this source further. Considering the uncertainties involved in

comparing our astrometry measurements with those made using other instruments, as noted above, the astrometry reported by Kouwenhoven et al. is in very good agreement with ours and further supports our conclusion that this source is co-moving with the primary star. Indeed, were the source a background star, its separation should have changed by  $-0''.1$  and its P.A. by  $-3'.3$  between 2000 and 2010, while the measured values show changes of  $+0''.01$  and  $+0'.6$ , respectively.

### 3.2. Properties of the Companion

In this section, in addition to comparing our spectrum and photometry measurements with those of other objects, we also rely on comparisons with synthetic spectra to estimate the



**Figure 2.** Measured (data points) separations (top) and position angles (bottom) between the primary star and its companion as a function of time. The solid lines show the expected position of a distant stationary background source over time, as calculated from the proper motion and distance of the primary star.

**Table 2**  
Properties of HIP 78530 AB

Parameter	Value	
	Primary	Companion
$\mu_\alpha \cos \delta$ (mas yr <sup>-1</sup> ) <sup>a</sup>	$-11.38 \pm 0.59$	...
$\mu_\delta$ (mas yr <sup>-1</sup> ) <sup>a</sup>	$-24.66 \pm 0.37$	...
Distance (pc) <sup>a</sup>	$156.7 \pm 13.0$	...
Angular separation (")	$4.533 \pm 0.006$	
Position angle (deg)	$140.3 \pm 0.1$	
$\Delta J$ (mag)	$8.13 \pm 0.05$	
$\Delta H$ (mag)	$7.44 \pm 0.03$	
$\Delta K'$ (mag)	$7.27 \pm 0.03$	
$\Delta K_{\text{cont}}^{2.09}$ (mag)	$7.26 \pm 0.03$	
$J$ (mag) <sup>b</sup>	$6.928 \pm 0.021$	$15.06 \pm 0.05$
$H$ (mag) <sup>b</sup>	$6.946 \pm 0.029$	$14.39 \pm 0.04$
$K_s$ (mag) <sup>b</sup>	$6.903 \pm 0.020$	$14.17 \pm 0.04$
$J - K_s$ (mag)	$0.025 \pm 0.03$	$0.89 \pm 0.06$
$H - K_s$ (mag)	$0.043 \pm 0.04$	$0.22 \pm 0.06$
Spectral type	B9V	M8±1
$T_{\text{eff}}$ (K)	$\sim 10500^c$	$2800 \pm 200$
$\log(L/L_\odot)$	...	$-2.55 \pm 0.13$
Mass ( $M_\odot$ )	$\sim 2.5^d$	$0.022 \pm 0.004^e$
Projected separation (AU)	$710 \pm 60$	

#### Notes.

<sup>a</sup> From van Leeuwen (2007).

<sup>b</sup> From the 2MASS PSC and our contrast measurements. Differences in filter bandpasses may add up to 1%–2% uncertainty.

<sup>c</sup> From the spectral type based on the temperature scale of Sherry et al. (2004).

<sup>d</sup> From the models of D'Antona & Mazzitelli (1997).

<sup>e</sup> From the models of Baraffe et al. (1998, 2002) and Burrows et al. (1997).

physical properties of the companion. For the comparison with models, we adopted the DRIFT-PHOENIX atmosphere models (Witte et al. 2009; Helling et al. 2008; Dehn et al. 2007), which combine a detailed kinetic model of dust cloud formation with a radiative transfer code (Hauschildt et al. 1999; Baron et al. 2003). In contrast to all other atmosphere models which assume phase equilibrium between gas and cloud particles, DRIFT-PHOENIX describes consistently the formation of a stationary cloud by homogeneous nucleation and grain growth/evaporation, including gravitational settling, element depletion, and convective element replenishment (Woitke & Helling 2003,

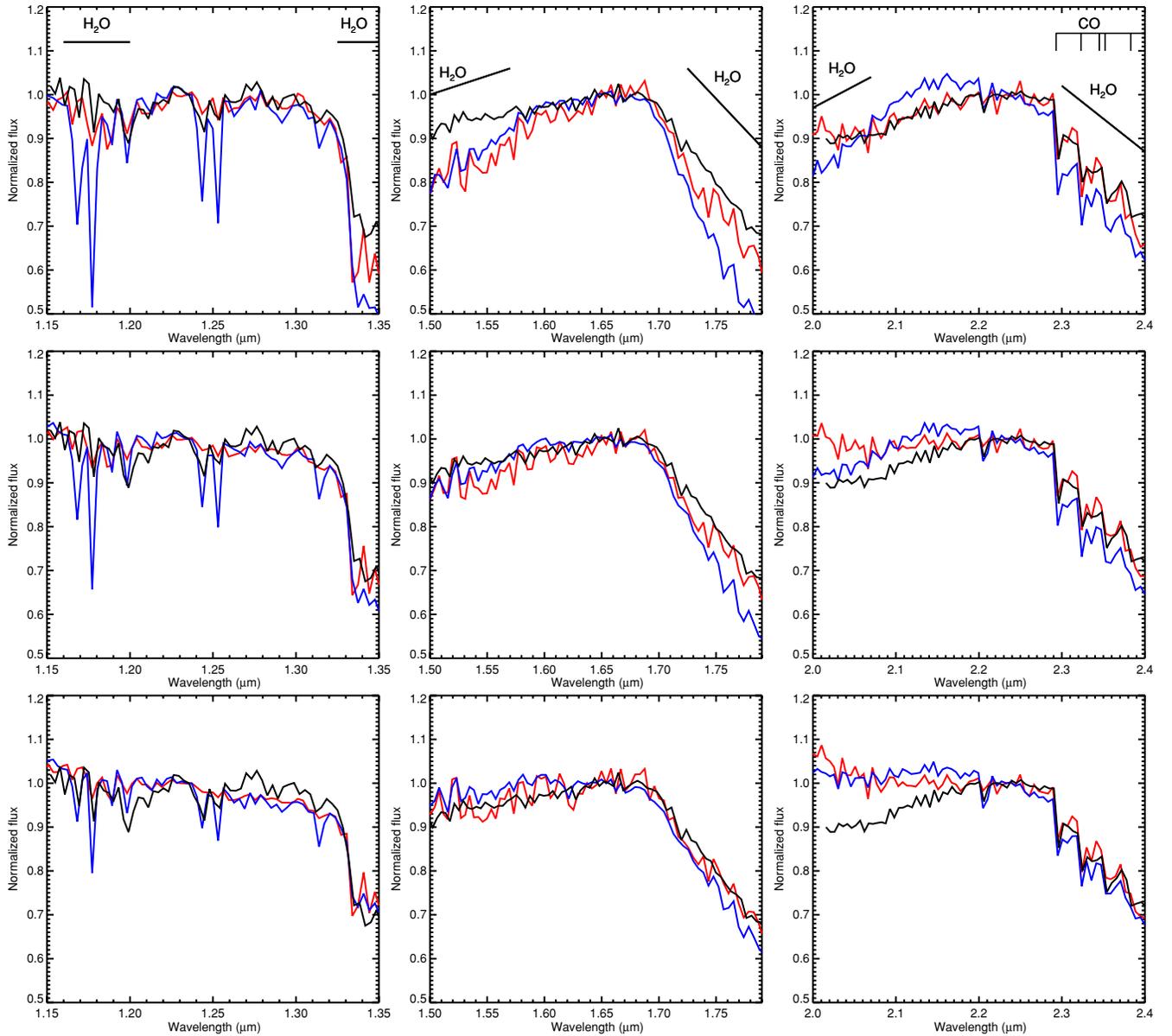
2004; Helling & Woitke 2006). All models used assume solar metallicity, which is reasonable for Upper Sco (Mohanty et al. 2004).

#### 3.2.1. Temperature, Surface Gravity, and Spectral Type

The spectrum of HIP 78530B, binned down to a resolving power of 400, is shown in Figure 3 in comparison with synthetic spectra from the DRIFT-PHOENIX atmosphere models for various temperatures and surface gravities. As shown in the figure, for all effective temperatures shown, the spectrum of the companion is in better agreement with the lower surface gravity spectrum, as expected for a young object in Upper Sco. This is particularly striking in the  $K$  band, where the broad  $\text{H}_2\text{O}$  absorption bands and the CO absorption band head depths are perfectly reproduced by the low gravity spectrum at  $T_{\text{eff}} = 2600$  K. The low surface gravity is also apparent in the depth of the  $K$  I lines in the  $J$  band. The effective temperature providing the best fit depends on the bandpass considered. As mentioned previously, a temperature of 2600 K provides an excellent fit in the  $K$  band. A temperature of 2600 K also provides a slightly better fit in the  $J$  band, both for the continuum around  $1.2 \mu\text{m}$  and the depth of the  $K$  I lines. In the  $H$  band, however, the best fit is obtained for temperatures of  $\sim 2800$ – $3000$  K. Taking these considerations into account, we adopt an effective temperature of  $2800 \pm 200$  K. We have also done a chi-square minimization using the same set of model spectra, rather than a simple visual inspection, and reached the same conclusion. In particular, when the minimization is done in all bands simultaneously, the best fit is found for a temperature of 2800 K.

Figure 4 shows the companion spectrum, at a higher resolution this time, in comparison with the spectrum of a known M8 dwarf member of Upper Sco as well as with a model spectrum with our effective temperature estimate of 2800 K. Among the Upper Sco brown dwarfs identified and spectrally classified by Lodieu et al. (2008), the best fit to our spectrum is obtained for a spectral type of M8, which corresponds to an effective temperature of  $\sim 2700$ – $2800$  K, consistent with our previous estimate. Comparing our spectrum to those of field dwarfs from the IRTF spectral library (Rayner et al. 2009) also results in a best fit for spectral types of M7–M9, although there are differences between the spectra owing to the lower gravity of the companion as noted earlier. The better agreement of our spectrum with those of other objects in Upper Sco, as opposed to field objects, provides further evidence that the new companion also belongs to the association.

The companion spectrum is not only in good agreement with those of models and other Upper Sco objects within the individual bands, but it also shows a reasonable agreement in colors between the bands. Over a temperature range of 2600–3000 K, the DRIFT-PHOENIX models yield synthetic colors of  $J - H = 0.50$ – $0.55$  and  $H - K = 0.27$ – $0.38$ , to be compared with the companion colors of  $J - H = 0.67 \pm 0.06$  and  $H - K = 0.22 \pm 0.06$ . Assuming a small possible extinction of  $A_V = 0.5$  mag toward this system (Carpenter et al. 2009), the extinction-corrected colors would be  $J - H \sim 0.62$  and  $H - K \sim 0.19$ . The colors of M8 dwarfs in Upper Sco from the sample of Lodieu et al. (2008) are 0.58–0.65 and 0.38–0.57, respectively. Here, the agreement in  $J - H$  is good but our companion is significantly bluer in  $H - K$ . We note that the differences between the NIRI (used for the companion contrast) and 2MASS (used for the primary star magnitude) filter bandpasses add an uncertainty of at most 2%–3% on the colors of the companions, calculated from synthetic spectra.



**Figure 3.** NIFS spectrum of HIP 78530B (black) compared with synthetic spectra from the DRIFT-PHOENIX atmosphere models of various effective temperatures and surface gravities. From the top to the bottom rows, the effective temperatures are 2600, 2800, and 3000 K, respectively. The red curves are for  $\log g = 4.0$  and the blue curves are for  $\log g = 6.0$ . The spectra were binned to a resolving power of 400 and were normalized separately in each spectral band.

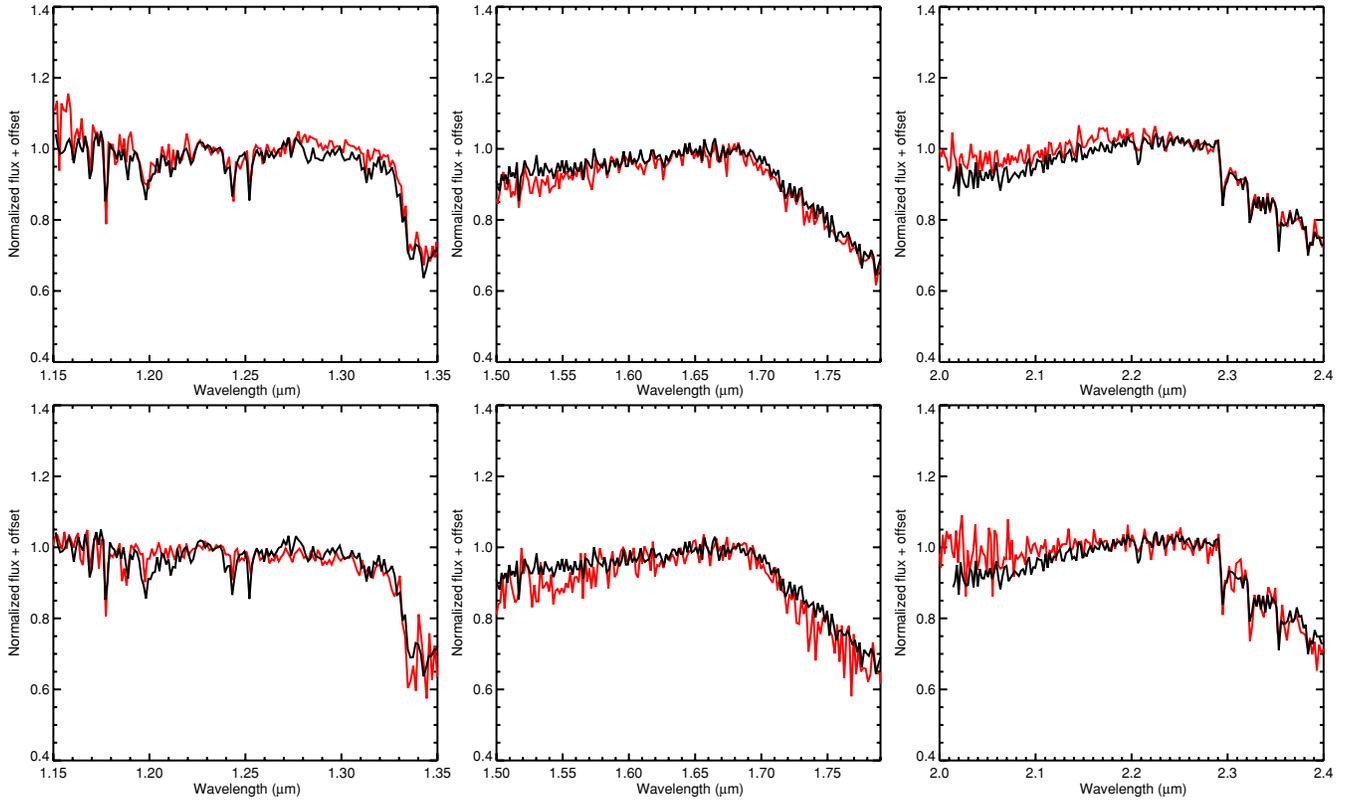
A good indicator of surface gravity is the equivalent width (EW) of the K I atomic line doublets at  $1.168/1.177 \mu\text{m}$  and  $1.243/1.253 \mu\text{m}$  (e.g., McGovern et al. 2004). For HIP 78530B, the EWs of these four lines are, respectively,  $2.0 \pm 0.3 \text{ \AA}$ ,  $3.0 \pm 0.3 \text{ \AA}$ ,  $2.0 \pm 0.2 \text{ \AA}$ , and  $2.1 \pm 0.2 \text{ \AA}$ .<sup>7</sup> In the DRIFT-PHOENIX models, there is a clear monotonic trend that the EW in these lines increases with increasing surface gravity for a constant temperature, see Figure 5. We have evaluated the EW for models of 2800 K with surface gravity ranging from 3.5 to 6.0 and interpolated the measured EW of the companion for each respective line. The mean and standard deviation of the four lines from this comparison indicate a surface gravity of  $\log g = 4.6 \pm 0.3$ , which is relatively low, as expected and verified for young objects (e.g., Brandeker et al. 2006). The K I

EWs we measured for the companion are also comparable to those of other M8–M9 members of Upper Sco (using spectra from Lodieu et al. 2008), but only a third to half of the EWs of M7–M9 field dwarfs of comparable spectral types (using spectra from McLean et al. 2003), clearly indicating low surface gravity and membership in Upper Sco for the new companion (see also Figure 5).

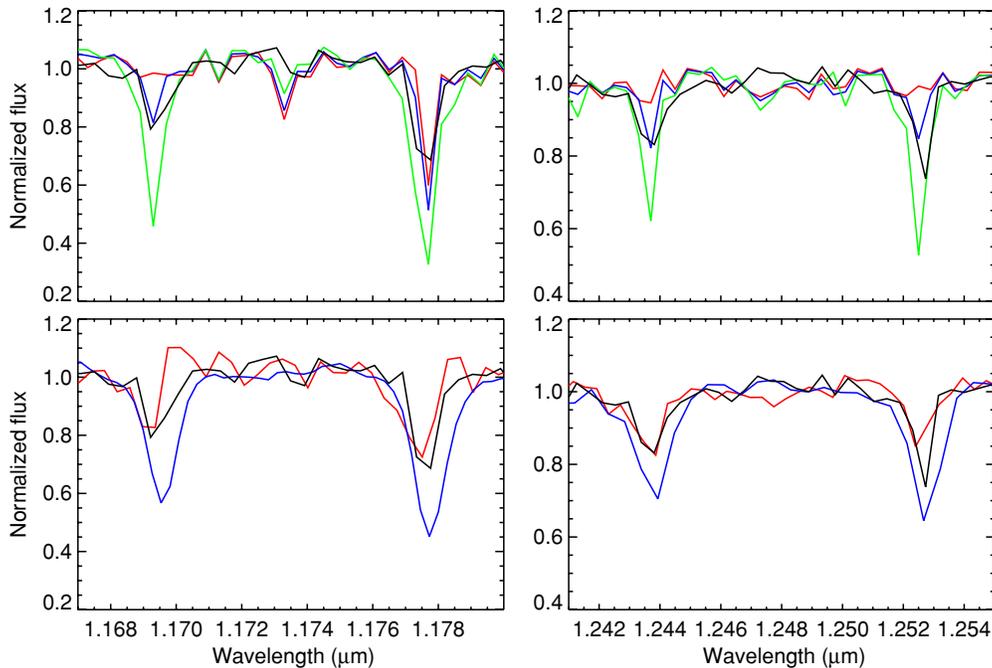
### 3.2.2. Atmosphere Composition and Structure

DRIFT-PHOENIX provides local gas properties such as the gas temperature  $T$  (K), the gas density  $\rho$  ( $\text{g cm}^{-3}$ ), and the local gas-phase composition, but also dust quantities such as the number of dust particles  $n_d$  ( $\text{cm}^{-3}$ ) of mean grain size  $\langle a \rangle$  (cm) at each layer of the atmosphere. These quantities are needed to evaluate the radiation transfer through the atmosphere taking into account convective energy transport, hence to calculate the synthetic spectrum as shown in Figure 3.

<sup>7</sup> To calculate the EWs, we used an integration interval of 10 nm centered on each line and fitted the continuum as a straight line using 40 nm wide intervals on both sides of the line.



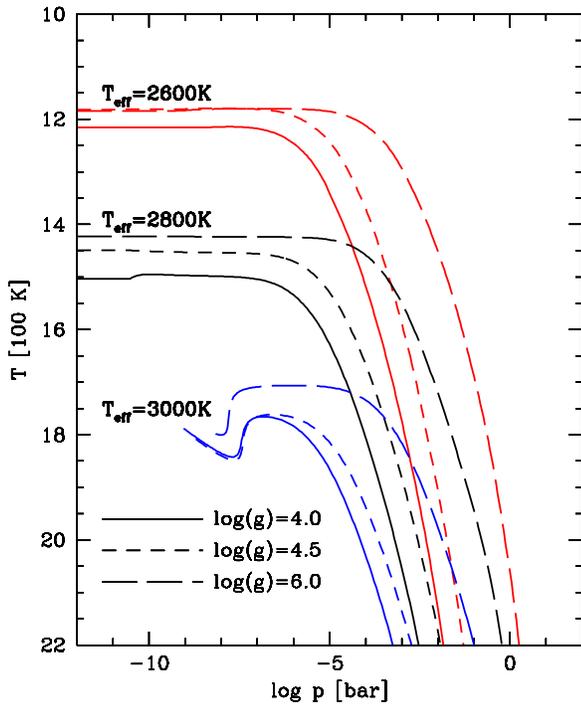
**Figure 4.** NIFS spectrum of HIP 78530B (black) compared with the spectrum of USco J155419.99–213543.1 (top), which is an M8 free-floating brown dwarf member of Upper Sco (Lodieu et al. 2008), and with a synthetic spectrum from the DRIFT-PHOENIX atmosphere models with  $T_{\text{eff}} = 2800$  K,  $\log g = 4.5$ , and solar metallicity (bottom). The spectra were binned to a resolving power of 1000 and were normalized separately in each spectral band.



**Figure 5.** Spectrum of HIP 78530B (black) showing the K I doublets at  $\sim 1.17$   $\mu\text{m}$  and  $\sim 1.25$   $\mu\text{m}$ . In the top panels, the companion spectrum is compared with spectra from the DRIFT-PHOENIX models for  $T_{\text{eff}} = 2800$  K and  $\log g = 3.5$  (red), 4.5 (blue), and 5.5 (green). In the bottom panels, the spectrum is compared with the spectrum of USco J155419.99–213543.1 (red, Lodieu et al. 2008), a young M8 brown dwarf in Upper Sco, and LP412-31 (blue, McLean et al. 2003), a field M8 dwarf. The shallower K I lines seen in the companion spectrum relative to the field object indicate lower surface gravity. The spectra were binned to a resolving power of  $\sim 3000$ .

The local gas-phase and cloud chemistry is determined by the local temperature and pressure which we demonstrate in Figure 6 for all DRIFT-PHOENIX atmosphere models applied in

Figure 3. While the low gravity spectra in the 2600–3000 K temperature range all provide a reasonable agreement with the observed spectrum, Figure 6 shows that the corresponding



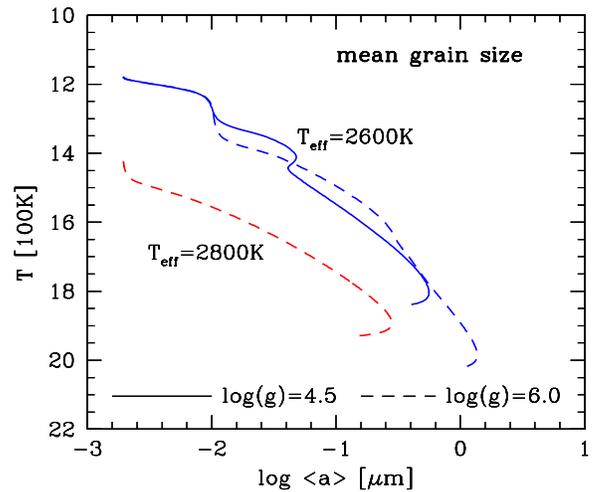
**Figure 6.** Temperature–pressure structures for all DRIFT-PHOENIX atmosphere models used in Figure 3. The  $T$ – $P$  profiles vary considerably for the three sets of models at different  $T_{\text{eff}}$ , although the emerging spectra from the three low-gravity atmosphere all provide a reasonable agreement with the observed companion spectrum.

(A color version of this figure is available in the online journal.)

underlying atmospheric ( $T$ ,  $p$ ) profiles are very different. This emphasizes that medium differences in the spectral energy distribution can hide large variation of the atmosphere’s structure. The models at 3000 K even develop a small temperature inversion high up in the atmosphere, where  $\text{H}_2\text{O}$  is dissociated, freeing up oxygen and locally enhancing the opacity compared to the surrounding layers.<sup>8</sup>

Of all the models shown in Figure 3, only three do have dust in their atmosphere:  $\log g = 4.5$  with  $T_{\text{eff}} = 2600$  K and  $\log g = 6.0$  with both  $T_{\text{eff}} = 2600$  K and  $2800$  K. Figure 7 shows the mean grain sizes for these models as a function of local temperature; the lowest temperatures indicate the upper part of the atmosphere. The dust persists at a higher temperature in the high-gravity models owing to an increased local density; however, these models are less relevant for the new companion found here, which is young and has low gravity. Thus, unless the companion is on the low end of our estimated  $T_{\text{eff}}$  range, it probably does not have dust in its atmosphere. However, if it is indeed closer to 2600 K and has dust in its atmosphere, the grain sizes would remain rather small, forming a haze layer in the upper, optically thin layers of the atmospheres (see Figure 7). The only objects where haze layers have been inferred from transit photometry, however, are irradiated giant gas planets (see Sing et al. 2009). No direct detection of haze on brown dwarfs has been obtained to date.

<sup>8</sup> Such a temperature and pressure inversion does not necessarily invalidate the assumption of hydrostatic equilibrium which is applied in the DRIFT-PHOENIX atmosphere models as was shown in Helling et al. (2000) and in Asplund (1998).



**Figure 7.** Mean grain sizes across the clouds altitudes in the three dust-forming model atmospheres plotted in Figure 3. The cloud grains remain rather small in all models and reside in the optical thin region of the atmosphere.

(A color version of this figure is available in the online journal.)

### 3.2.3. Luminosity and Mass

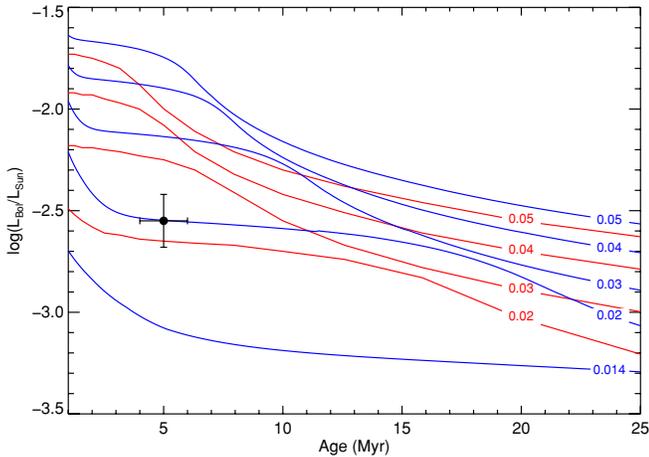
We estimated the bolometric luminosity of the companion by using its observed photometry in combination with model spectra. We first computed synthetic fluxes in the  $JHK$  bands using the model spectra, then adjusted those synthetic fluxes to the measured values, and finally integrated the scaled model spectra over all wavelengths to obtain the total irradiance, which was then converted to bolometric luminosity using the primary star distance. The synthetic average flux densities in the  $JHK$  bands were computed using the relative spectral response curves given on the Two Micron All Sky Survey (2MASS) project Web site.<sup>9</sup> We repeated the procedure using both the DRIFT-PHOENIX and NextGen (Hauschildt et al. 1999) model spectra for ranges of  $T_{\text{eff}}$  from 2600 K to 3000 K and  $\log g$  from 3.5 to 5.0. We obtained a value of  $\log(L/L_{\odot}) = -2.55 \pm 0.13$ . The uncertainty quoted reflects the range of values obtained for the different model parameters and accounts for the uncertainty on the distance of the primary star.

The age of the Upper Scorpius association is well constrained at 5 Myr and all stars appear to have formed in a burst, over a period of at most 1–2 Myr (Preibisch et al. 2002; Slesnick et al. 2008). For an age of  $5 \pm 1$  Myr and the above companion bolometric luminosity estimate, the evolution models of Chabrier et al. (2000) indicate a mass of  $\sim 21$ – $26 M_{\text{Jup}}$ , while those of Burrows et al. (1997) yield a mass of  $\sim 19$ – $25 M_{\text{Jup}}$ , see Figure 8. The ranges of mass quoted reflect only the uncertainties on the age and luminosity estimates, but of course, larger uncertainty may be present owing to the lack of absolute calibration of the evolution models. As visible in Figure 8, the companion is currently in a phase where its luminosity is more or less constant with time due to deuterium burning; this phase will last until an age of  $\sim 15$  Myr. Interestingly, this means that the companion mass estimate is not strongly dependent on the age estimate, as is ordinarily the case for substellar objects.

## 4. DISCUSSION AND CONCLUDING REMARKS

Due to the very large semimajor axis of the companion orbit, its expected orbital motion per year is below the astrometric

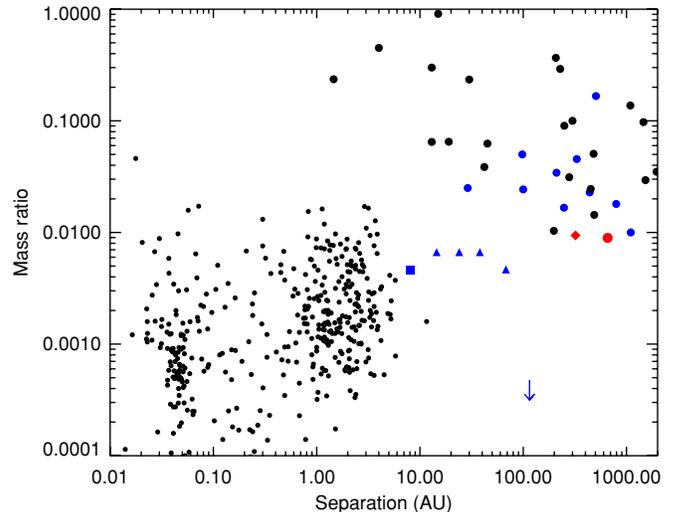
<sup>9</sup> [http://www.ipac.caltech.edu/2mass/releases/allsky/doc/sec6\\_4a.html](http://www.ipac.caltech.edu/2mass/releases/allsky/doc/sec6_4a.html)



**Figure 8.** Calculated bolometric luminosity of the companion (black point) compared with evolutionary tracks from Baraffe et al. (1998, 2002) (red curves) and Burrows et al. (1997) (blue curves). The curves are labeled with their mass expressed in units of solar mass.

precision of our data. To estimate the orbital motion, we can consider the hypothetical case in which the orbit is circular and face-on as seen from Earth, with a semimajor axis of 710 AU. Given the system mass of approximately  $2.5 M_{\odot}$ , this yields an orbital period of  $\sim 12,000$  years, which in turn gives an angular motion of  $0.03 \text{ yr}^{-1}$  (or a linear motion of  $\sim 2.4 \text{ mas yr}^{-1}$ ). The angular motion over a decade under this assumption,  $0.3$ , does come out to the same order of magnitude as the  $2\sigma$  difference in P.A. between our images and those of Kouwenhoven et al. (2005),  $0.6 \pm 0.3$ . However, in reality, the orbital motion is likely to be much smaller, as the projected separation is typically smaller than the real semimajor axis. In addition, the only way in which the angular motion could be larger than  $0.3$  is if the orbit is eccentric and the companion is presently close to periastron, which is a priori unlikely. As we note in Section 3.1, our astrometry is internally consistent for the NIRI data but not necessarily with respect to other instruments, hence it is plausible that the difference with Kouwenhoven et al. is simply spurious. Thus, there is no convincing evidence of orbital motion, and no such motion is expected over the relevant timescales.

A mass ratio-separation diagram showing HIP 78530B and other known brown dwarf and planetary companions to stars is shown in Figure 9. With a mass of  $\sim 23 M_{\text{Jup}}$  and projected separation of  $\sim 700$  AU from its massive  $2.5 M_{\odot}$  primary, HIP 78530B lies at the lower boundary of all known companions with separations larger than 100 AU. It is unclear whether these companions, whose mass ratios overlap with those of both more massive brown dwarf companions and bona fide planets, formed in a planet-like or in a stellar-like manner. Given the size of their orbits, it appears unlikely that they formed in situ in a planet-like manner—either through the collapse of a gravitationally unstable disk (e.g., Vorobyov & Basu 2010) or by the core accretion mechanism (e.g., Pollack et al. 1996). Indeed for both scenarios, the disks would need to be unusually large and massive, but even then, the formation timescale by core accretion would be prohibitively long while the efficiency of disk instability to produce such companions is highly uncertain (Vorobyov & Basu 2010; Kratter et al. 2010). Nevertheless, it is possible that these companions did form in a planet-like manner much closer to the star, but were subsequently kicked outward through gravitational interactions



**Figure 9.** Mass ratio as a function of separation for various substellar companions to stars. The red filled circle is HIP 78530B and the red filled diamond is 1RXS J160929.1–210524b. The blue filled circles are all other directly imaged low-mass substellar companions ( $< 25 M_{\text{Jup}}$ ) to stars from the compilation given in Lafrenière et al. (2010), augmented with the newly found companion Ross 458C (Goldman et al. 2010; Scholz 2010). The large black filled circles are more massive directly imaged substellar companions from the compilation of Zuckerman & Song (2009). The triangles, square, and down arrow are HR 8799bcde (Marois et al. 2008, 2010),  $\beta$  Pic b (Lagrange et al. 2009, 2010), and Fomalhaut b (Kalas et al. 2008), respectively. The small black circles indicate planets found by the radial velocity and microlensing techniques from the Extrasolar Planets Encyclopaedia (<http://exoplanet.eu/>).

(e.g., Scharf & Menou 2009; Veras et al. 2009). In this scenario, according to the simulations of Veras et al. (2009), the timescale for instabilities to develop and send planets on large orbits may be quite short (0.01–1 Myr), with significant dynamical evolution occurring within the first few Myr. However, owing to further evolution, planets quickly scattered on large orbits are likely to be ejected from the system after a few tens of Myr. Alternatively, these wide and low-mass ratio companions could form like stellar binaries, through the fragmentation of a pre-stellar core (e.g., Bate 2009; Bate et al. 2003). Numerical simulations indicate that very low mass ratio companions can indeed be produced by this process, albeit rarely, and that they preferentially have large separations and are generally found in high-order multiple systems. Future observations may provide constraints to exclude or support the various formation possibilities of these wide low-mass companions. For instance, if they formed in a planet-like manner closer to the star and were subsequently kicked outward through gravitational interactions, then it would be expected that additional companions of similar mass or even heavier should be present in the systems, at smaller separations.

In addition, as mentioned above, distant companions produced through this process would likely be eventually ejected from the system, such that they should be found more frequently in star-forming regions than around older stars. In our overall survey of Upper Scorpius that led to the discovery of HIP 78530B and 1RXS J1609–2105b, we observed 91 stars (masses  $0.15\text{--}5 M_{\odot}$ ), so taken at face value, our two discoveries imply that companions with mass ratios below 0.01 and separations of hundreds of AUs exist in  $2.2^{+5.5}_{-1.9}\%$  (95% credibility) of stellar systems. Considering that we did not achieve the same sensitivity for all targets as well as our incompleteness to lower mass ratio companions, this number is only a lower

limit. The statistics are not yet sufficient to tell whether there exists a difference for older objects. For example, the study of Lafrenière et al. (2007), which targeted  $\sim 200$  Myr-old GKM stars, enabled placing upper limits of  $\sim 6\%$  (95% credibility) for companions of  $\sim 10 M_{\text{Jup}}$  (i.e., mass ratio  $\sim 0.01$ ); see also Chauvin et al. (2010) and Nielsen & Close (2010). Improving the statistics for both younger and older systems would thus be a good way to investigate further the importance of gravitational scattering in accounting for distant companions. This approach is valid as long as the internal dynamics of the multiple planets dominates over interactions with other stars or giant molecular clouds as the system travels through the galaxy, as these interactions can also strip out wide companions. This can be roughly checked using the results of Weinberg et al. (1987). Based on their Figure 2, with  $a/M_{\text{tot}} \sim 0.0014 \text{ pc } M_{\odot}^{-1}$ , the disruption lifetime of HIP 78530AB due to encounters with stars and giant molecular clouds is larger than 10 Gyr; this is also the case for 1RXS J160929.1–210524Ab. In other words, stellar and molecular clouds encounters have little impact on the evolution of companions such as the ones we have found.

The new companion reported in this paper joins a growing list of low-mass substellar companions ( $\lesssim 25 M_{\text{Jup}}$ ) in wide orbit ( $\gtrsim 100$  AU) around stars, which now counts about 10 objects (see Lafrenière et al. 2010 for a recent compilation, see also Figure 9). These companions are found around primaries covering a wide range of masses and even around primaries that are themselves binaries. Thus, while the statistics on their frequency may not be very accurate, it seems likely that wide low-mass substellar companions are not an unusual outcome of the star formation process, yet they remain hard to explain within current theoretical frameworks. Over the next few years, additional searches for new companions and continued efforts to characterize the known ones should allow us to make good progress toward explaining their formation.

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